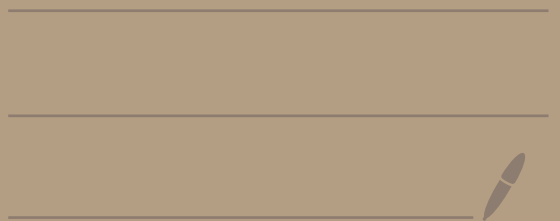


# Stability of Matter p.3

5. Stability of 1st kind



## 5. Stability of first kind

Notation: we will say that  $f \in L^p + L^q$  if  $f = f_1 + f_2$  where  $f_1 \in L^p, f_2 \in L^q$ .

Equipped with the Sobolev inequality we are ready to prove

### Thm (stability of first kind, $d=3$ )

Let  $d=3$  and assume that

$$H = -\Delta + V \quad \text{with} \quad V \in L^{\frac{3}{2}}(\mathbb{R}^3) + L^\infty(\mathbb{R}^3)$$

Then  $E_0 > -\infty$ .

$$V_- = \max\{0, -V\}$$

Proof.

Recall the Sobolev inequality in  $d=3$ :

$$\|\nabla f\|_{L^2}^2 \geq C_3 \|f\|_6^2$$

with  $C_3$  a constant. The optimal constant is in fact  $C_3 = \frac{2}{9} (4\pi^2)^{2/3}$ .

Recall

$$E(\psi) = \frac{1}{2} \int_{\mathbb{R}^3} |\nabla \psi|^2 dx + \int_{\mathbb{R}^3} V(x) |\psi(x)|^2 dx. \quad \text{Thus}$$

$$E(\psi) \geq \frac{C_3}{2} \|\psi\|_6^2 + \int_{\mathbb{R}^3} V(x) |\psi(x)|^2 dx \geq \frac{C_3}{2} \|\psi\|_6^2 - \int_{\mathbb{R}^3} V_-(x) |\psi(x)|^2 dx$$

We could now use Hölder's inequality

$$\|fg\|_1 \leq \|f\|_p \|g\|_q \quad \frac{1}{p} + \frac{1}{q} = 1$$

to obtain

$$\int_{\mathbb{R}^3} V_-(x) |\psi(x)|^2 \leq \|V_-\|_{\frac{3}{2}} \|\psi\|_6^2$$

then

$$E(\psi) \geq \frac{C_3}{2} \|\psi\|_6^2 - \|V_-\|_{\frac{3}{2}} \|\psi\|_6^2 \geq 0$$

when  $\|V_-\|_{\frac{3}{2}} \leq \frac{C_3}{2}$ . (\*\*)

Of course  $\frac{1}{|x|}$  is not in  $L^{\frac{3}{2}}$ . This is why we need the  $L^\infty$  part. To this end we claim that there exists a constant  $\lambda < 0$  such that

$$h(x) = \min(V(x) - \lambda, 0) \leq 0$$

satisfies  $\|h\|_{\frac{3}{2}} \leq \frac{C_2}{2}$ . (\*) (proved below  $\rightarrow$  exercise)

Let  $V(x) = v_1 + v_2$  with  $v_1 \in L^{\frac{3}{2}}$ ,  $v_2 \in L^\infty$ .

We have

$$E(\psi) = T_\psi + \int_{\mathbb{R}^3} (v_1(x) - \lambda) |\psi(x)|^2 dx + \lambda \int_{\mathbb{R}^3} |\psi(x)|^2 dx$$

$\geq 0$  by (\*\*)

$$+ \int_{\mathbb{R}^3} v_2(x) |\psi(x)|^2 dx \geq \frac{1}{2} (T_\psi + 2 \langle v_1, h \psi \rangle)$$

$$+ \frac{1}{2} T_\psi + \lambda - \|v_2\|_\infty \geq \frac{1}{2} T_\psi + \lambda - \|v_2\|_\infty > -\infty$$

## Exercise

$$h(x) = \min(V(x) - \lambda, 0), \quad V(x) \leq 0, V \in L^{\frac{3}{2}}.$$

Then  $\forall \varepsilon > 0$   $\exists$  exists a  $\lambda$  such that  $\|h\|_{\frac{3}{2}} < \varepsilon$ .

Solution:

$$\|h\|_{\frac{3}{2}}^{\frac{3}{2}} = \int_{\mathbb{R}^3} \overset{h \leq 0}{(-\min(V(x) - \lambda, 0))}^{\frac{3}{2}} dx = \int_{V \leq \lambda} (|V(x)| - |\lambda|)^{\frac{3}{2}}$$

$$\leq \int_{|V| \geq |\lambda|} |V(x)|^{\frac{3}{2}} dx \longrightarrow 0 \quad \text{as } |\lambda| \rightarrow \infty$$

since  $V \in L^{\frac{3}{2}}$

(indeed by monotone conv.  $\int |f| \chi_{|f| \leq M} \xrightarrow{M \rightarrow \infty} \int |f|$ )

What about other dimensions?

The Sobolev inequality that we have proved works in  $d \geq 2s$  where  $s=1$  in the non-relativistic case. Thus we can consider  $d \geq 3$ , too.

## Exercise

Derive stability of first kind for non-relativistic matter in  $d \geq 3$ .

Solution

By the Sobolev inequality for  $d \geq 3$

$$|\nabla \varphi|_2^2 \geq C \|\varphi\|_{\frac{2d}{d-2}}^2.$$

In particular we have

$$\int_{\mathbb{R}^d} |V_-(x)| |\psi(x)|^2 dx \leq \left( \int_{\mathbb{R}^d} |V_-|^{\frac{d}{2}} \right)^{2/d} \left( \int_{\mathbb{R}^d} |\psi|^{\frac{2d}{d-2}} \right)^{\frac{d-2}{d}}$$

$$\frac{2}{d} + \frac{1}{q} = 1 \quad \frac{1}{q} = 1 - \frac{2}{d} = \frac{d-2}{d} \Rightarrow q = \frac{d}{d-2}$$

We can now repeat all steps from  $d=3$  case to conclude that

$$E_0 > -\infty \quad \text{if} \quad V \in L^{\frac{d}{2}}(\mathbb{R}^d) + L^\infty(\mathbb{R}^d), \quad d \geq 3$$

### Lower dimensions $d=1,2$

We now ask the same question for  $d=1,2$ .

### Thm (Sobolev's inequality in $d=1,2$ )

a) For any  $f \in H^1(\mathbb{R})$  we have

$$\|f\|_\infty^2 \leq \|f'\|_2 \|f\|_2$$

Furthermore

$$|f(x) - f(y)| \leq \|f'\|_2 |x-y|^{1/2}.$$

(b) For any  $f \in H^1(\mathbb{R}^2)$  we have

$$C_{2,p} \|f\|_p^2 \leq \|\nabla f\|_2^2 + \|f\|_2^2 \quad \forall p \in [2, \infty)$$

## Sketch of proof

(a) Assume  $f \in C_0^\infty(\mathbb{R})$ . Then

$$f(x)^2 = \int_{-\infty}^x f(y) f(y)' dy - \int_x^{\infty} f(y) f(y)' dy$$

Thus

$$|f(x)|^2 \leq \int_{-\infty}^x |f| |f'| + \int_x^{\infty} |f| |f'| = \int_{-\infty}^{\infty} |f| |f'|$$

$$\stackrel{CS}{\Rightarrow} \|f\|_\infty^2 \leq \|f'\|_2 \|f\|_2$$

One then uses a density argument since  $C_0^\infty$  is dense in  $H^1$ .  
The other bound follows from the fact that

$$|f(x) - f(y)| = \left| \int_y^x f'(t) dt \right| \leq \left( \int_y^x 1 dt \right)^{1/2} \left( \int_y^x |f'(t)|^2 dt \right)^{1/2} \leq |x-y|^{1/2} \|f'\|_2$$

this follows since if  $f \in H^1(\mathbb{R}) \Rightarrow f(x) = c + \int_a^x f'(t) dt$  a.e.

(obvious for smooth fcts).

b) Let  $q \in (1, 2)$ . Then

$$\begin{aligned} \|\vec{f}\|_q^q &= \int_{\mathbb{R}^2} |\vec{f}(u)|^q du = \int_{\mathbb{R}^2} \frac{|\vec{f}(u)|^q (1 + 4\pi^2 |u|^2)^{\frac{1}{2}}}{(1 + 4\pi^2 |u|^2)^{\frac{1}{2}}} du \\ &\leq \underbrace{\left( \int_{\mathbb{R}^2} |\vec{f}(u)|^2 (1 + 4\pi^2 |u|^2) du \right)^{\frac{1}{2}}}_{H^1 \text{ norm}^2} \left( \int_{\mathbb{R}^2} \frac{du}{(1 + 4\pi^2 |u|^2)^{\frac{1}{2}}} \right)^{1 - \frac{q}{2}} \end{aligned}$$

where  $\frac{1}{2-q} + \frac{1}{2} = 1$   $h = \frac{q}{2} \cdot z = \frac{q}{2} \frac{2}{2-q} = \frac{q}{2-q} > 1$

$$\frac{1}{2} = 1 - \frac{q}{2} \Rightarrow z = \frac{2}{2-q}$$

This implies  $\int_{\mathbb{R}^2} \frac{dx}{(1 + \frac{q}{2} |x|^2)^2} dx < \infty$

$$\Rightarrow \|f\|_q^q \leq C \|f\|_{H^1}^q \Rightarrow \|f\|_q \leq C \|f\|_{H^1}$$

Using the Hausdorff-Young inequality  $\|f\|_p \leq \|f\|_q$  for  $\frac{1}{p} + \frac{1}{q} = 1, q \in [1, 2]$ . Since  $1 < q < 2 \Rightarrow p \in (2, \infty)$ .

For  $p=2$  it is obviously true. ■

## Exercise

Show the following theorem:

Thm (Stability of matter in lower dimensions)

Let  $V \in L^1(\mathbb{R}) + L^\infty(\mathbb{R})$  in  $d=1$  and

$V \in L^{1+\epsilon}(\mathbb{R}^d) + L^\infty(\mathbb{R}^d)$  in  $d \geq 2$ . then

$$E_0 > -\infty.$$

## Proof / solution

We follow the argument for  $d \geq 3$ . This time we have

$$\underline{d=1} \quad \int_{\mathbb{R}} |V(x)| |\psi(x)|^2 dx \leq \|V\|_1 \|\psi\|_\infty^2$$

$$\text{and } T_\psi \geq C \|\psi(x)\|_\infty^4$$

$$\Rightarrow E(\psi) \geq C \|\psi(x)\|_\infty^4 - \|V\|_1 \|\psi\|_\infty^2 > -\infty$$

Note that here using the  $L^\infty$  is not needed but allows for more potentials.

$d=2$ .

$$\text{Since } \|\nabla f\|_2^2 \geq C_{hp} \|f\|_p^2 - \|f\|_2^2$$

$$\Rightarrow E(\psi) \geq C \|\psi\|_p^2 - C - \int_{\mathbb{R}^2} |V(x)| |\psi(x)|^2 dx$$

if  $V \in L^\infty \Rightarrow$  the  $p=2$  case done

if  $V \in L^{1+\varepsilon}$

$$\int_{\mathbb{R}^2} |V| |\psi|^2 \leq \left( \int_{\mathbb{R}^2} V^{1+\varepsilon} \right)^{\frac{1}{1+\varepsilon}} \left( \int_{\mathbb{R}^2} |\psi|^{2\tilde{p}} \right)^{\frac{1}{\tilde{p}}}$$

$$\frac{1}{1+\varepsilon} + \frac{1}{\tilde{p}} = 1 \Rightarrow \frac{1}{\tilde{p}} = 1 - \frac{1}{1+\varepsilon} = \frac{\varepsilon}{1+\varepsilon}$$

$$\Rightarrow \tilde{p} = \frac{1+\varepsilon}{\varepsilon} \rightarrow \infty \text{ as } \varepsilon \rightarrow 0$$

so for  $\forall \varepsilon > 0$   $\exists \tilde{p} < \infty$  :

$$2\tilde{p} = q \quad \frac{1}{\tilde{p}} = \frac{2}{q}$$

$$E(\psi) \geq C \|\psi\|_{\tilde{p}}^2 - C - \tilde{C} \|\psi\|_{\tilde{p}}^{\frac{2}{q}}$$

we use the  $2$ -trick  $L^\infty$  to get the bound  
since then we can make  $\tilde{C} < C$ . ■